





On the search for the most energetic <u>photons</u> in the <u>Universe</u>

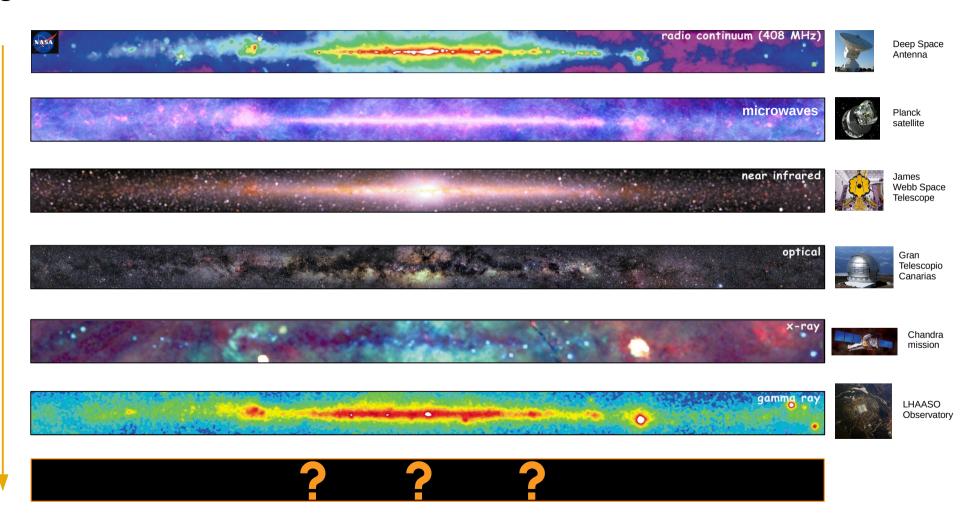
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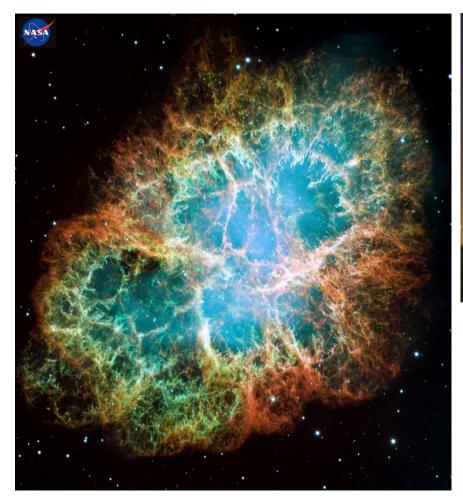


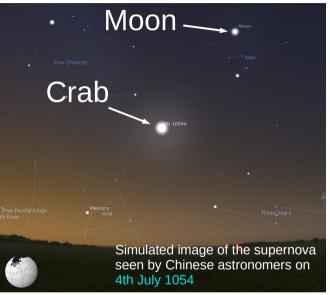


Light comes in different flavours



Crab Nebula, the most studied gamma-ray source





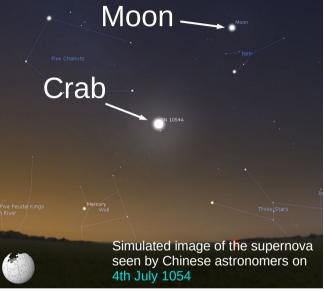


Holy Roman Emperor Henry III before Tivoli pointing up at a new star (Univ. Heidelberg)

- ▶ Remnants of a star explosion leaving behind a rotating neutron star
- ▶ 42 gamma-rays per second hitting an equivalent area of Buenos Aires
- ► Steep energy spectrum: the higher the energy of the gamma-ray, the more difficult to produce it

TeV photons from Crab (1989)



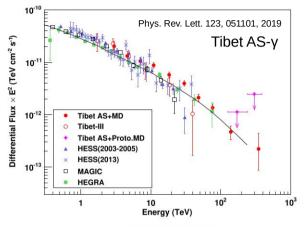


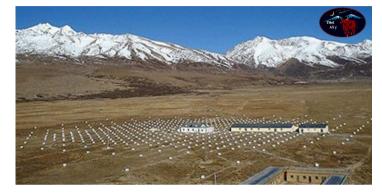


Holy Roman Emperor Henry III before Tivoli pointing up at a new star (Univ. Heidelberg)

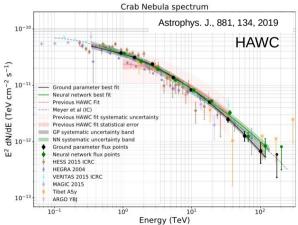
- ▶ Remnants of a star explosion leaving behind a rotating neutron star
- ▶ 42 gamma-rays per second hitting an equivalent area of Buenos Aires
- ► Steep energy spectrum: the higher the energy of the gamma-ray, the more difficult to produce it
- ► TeV gamma-ray source discovery by Whipple, an imaging Cherenkov telescope with a 10-m reflector and a multi-pixel camera, in 1989 (Astrophys. J. 342 (1989) 379)

100 TeV photons from Crab (2019)





4300 m above sea level (Yangbajing, Tibet)





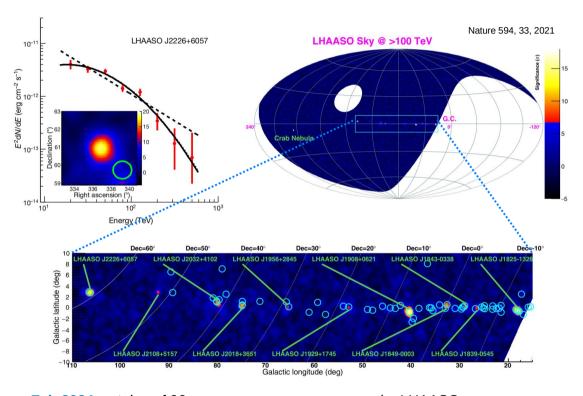
4100 m above sea level (Puebla, Mexico)

1 PeV photons from Crab! (2021)

4400 m above sea level (Sichuan, China)

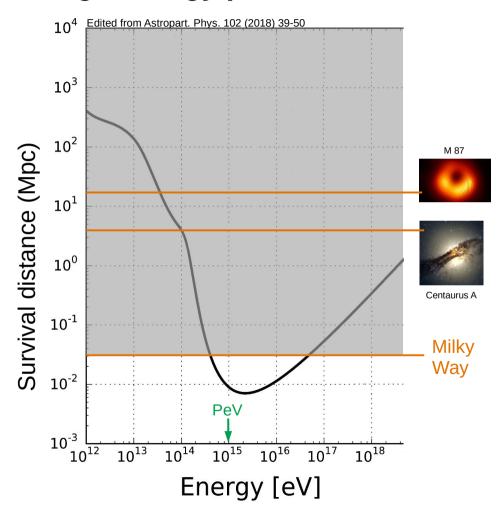


May 2021: 12 sources of gamma-rays with energies up to 1 PeV



Feb 2024: catalog of 90 gamma-ray sources sources by LHAASO (ApJS 271, 25, 2024)

Ultra-high-energy photons



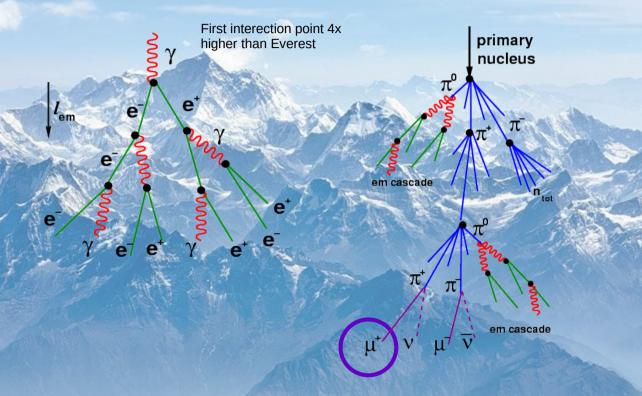
Gamma-rays with energies above few PeV are usually called "ultra-high-energy photons" (UHE)

- ► Between 1-100 PeV, photons trace the local Universe → messengers of Galactic phenomena
- ► Produced in cosmic-ray acceleration sites, during cosmic-ray propagation or in the decay of dark matter
- ► They carry information about:
 - → high-energy particle interactions
 - → Galactic matter and radiation fields
 - → beyond standard model Physics

► Their detection is not an easy task: less than 1 UHE photon per year hitting an equivalent area of Buenos Aires (E > 50 PeV)

Then, how could we study them?

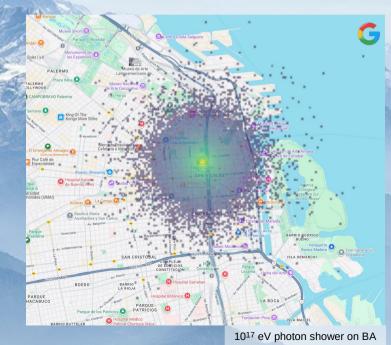
Is this a lottery for ultra-high-energy photons?



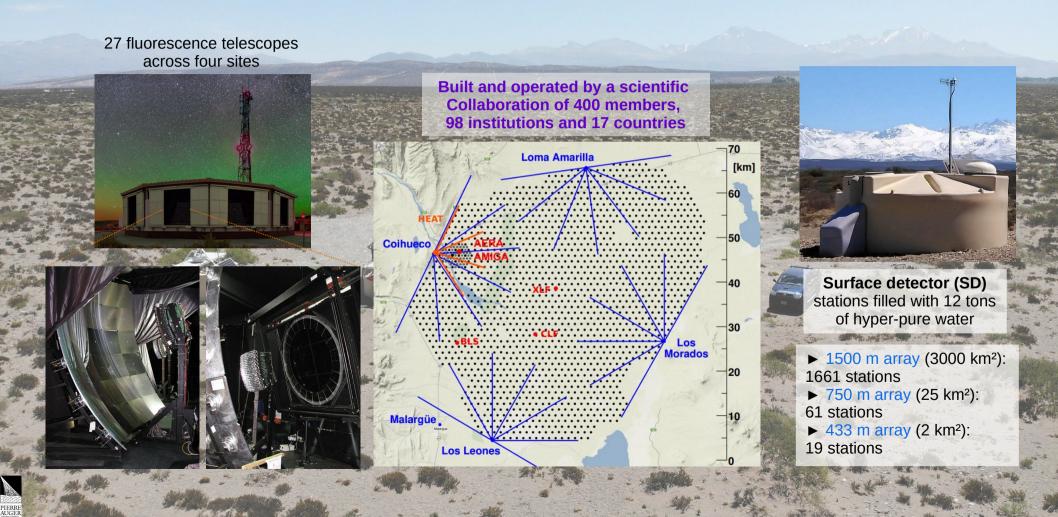
- ► Cosmic rays outnumber photons in at least 1:100,000 (same odds to winning the roulette three times in a row)
- ► Cosmic rays produce hadronic cascades with a rich muon component

We can discover ultra-high-energy photons by measuring air-shower muons

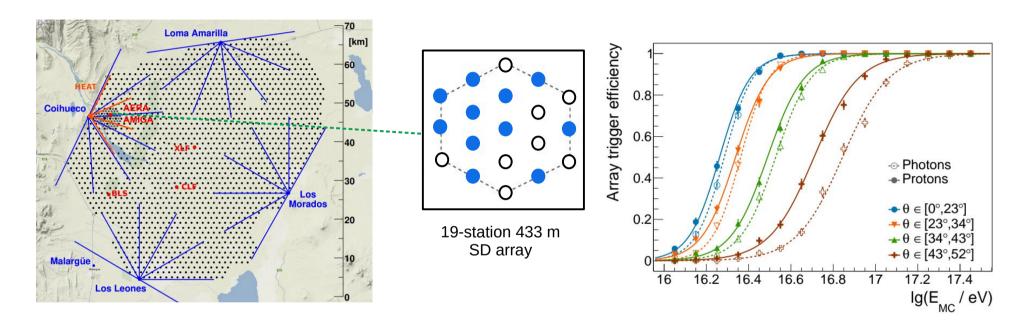
► Astroparticles at ultra-high energies interact with the atmosphere initiating a cascade of secondary leptons and hadrons



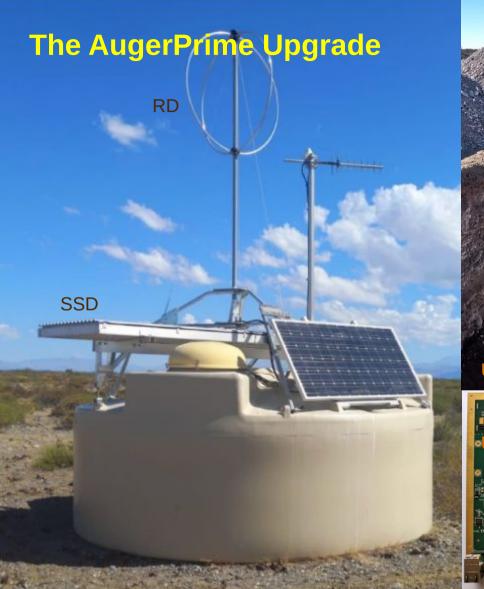
The Pierre Auger Observatory



Cosmic-ray detection above 30 PeV with the SD-433



- ► Significant air-shower detection probability above 30 PeV → we can infer the features of the primary particle with the SD-433
 - ► Direct muon measurement with one of the detectors of the AugerPrime upgrade: the Underground Muon Detector (UMD)





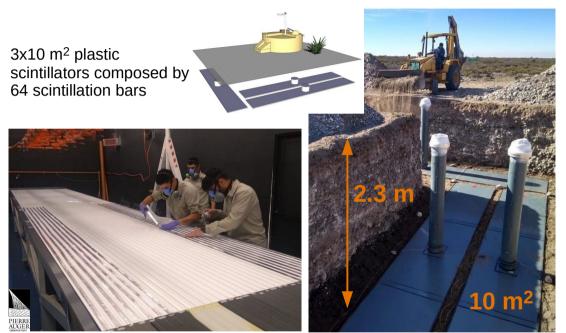


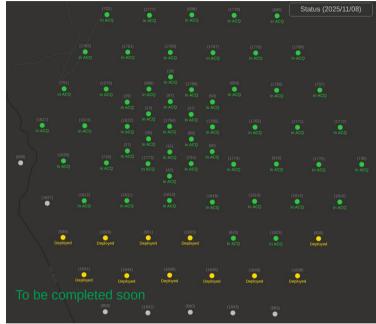
Multi-hybrid air-shower detection:

- ► 433 and 750 m SD arrays equipped with 3 x 10 m² UMD stations
- ► 4 m² plastic surface scintillation detectors (SSD)
- ► Measure radio emission from air showers with antennas (RD)
- ► Faster electronics (UUB) and larger dynamic range thanks to a small-area PMT (SPMT)

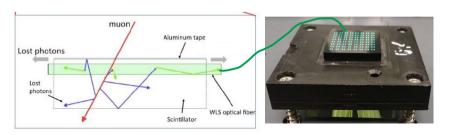
Aiming at a better composition sensitivity (including UHE γ)

Using buried scintillators to win the lottery





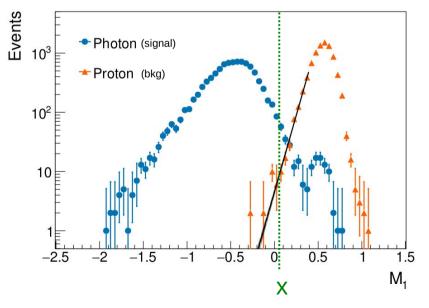
- ► All charged particles from the air-shower are filtered in the ground, except muons with kinetic energies above 1 GeV
- ► Plastic scintillation material emits fluorescence light when a muon transverses it
- ► Single-pixel silicon photomultipliers (SiPM) as optosensors



How we searched for photons at tens of PeV

- \blacktriangleright Discrimination based on the muon density (ρ_i) measured by the UMD stations
- ► The normalization factor ρ_{pr} is the expected muon density for proton events at r_{pr} = 200 m

$$M_1 = \lg \left(\sum_i \frac{\rho_i}{\rho_{\rm pr}} \times \left(\frac{r_i}{r_{\rm pr}} \right) \right)$$



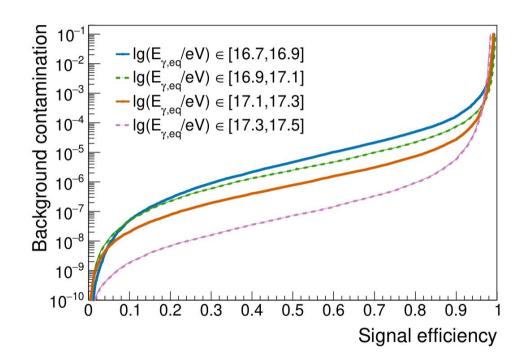
► Two metrics characterize the discrimination performance:

sig. efficiency =
$$\int_{-\infty}^{x} f_{\gamma}(M_1) dM_1$$

bkg. contamination =
$$\int_{-\infty}^{x} f_{pr}(M_1) dM_1$$

M₁: the discrimination observable

 \blacktriangleright Discrimination based on the muon density (ρ_i) measured by the UMD stations



► Two metrics characterize the discrimination performance:

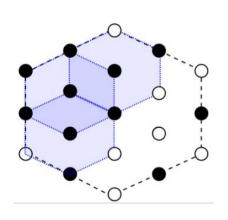
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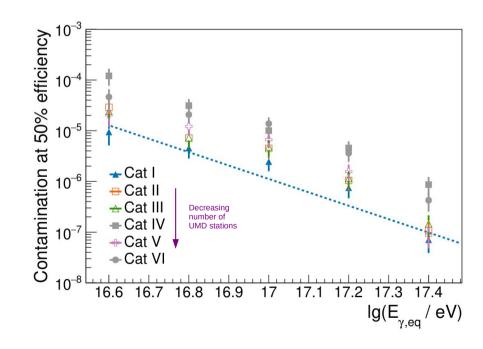
- ► Contamination decreases with increasing primary energy due to the larger air-shower muon content
- ▶ Larger signal efficiency at the expense of larger contamination

Background contamination

- ▶ M₁ calculated with the UMD stations of the hottest hexagon
- ▶ Since M₁ scales with the number of muon density measurements, any missing UMD station would bias M₁
 - → event classification in six "categories" based on number and placement of available UMD stations



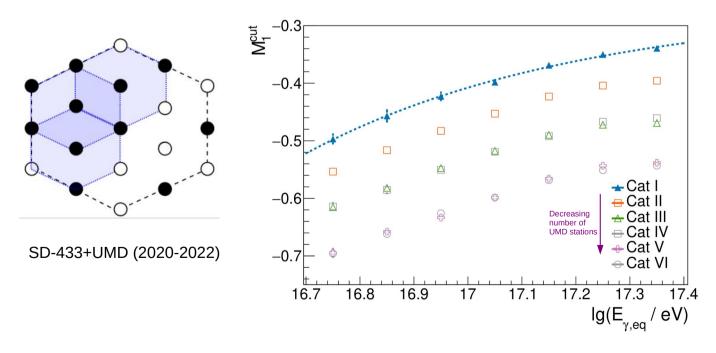
SD-433+UMD (2020-2022)



- ► Contamination increases for decreasing number of UMD stations
- ► Very low contamination!
- 1 bkg event misidentified as a photon every 106 years in Cat I (NW hexagon)

Photon candidate cuts

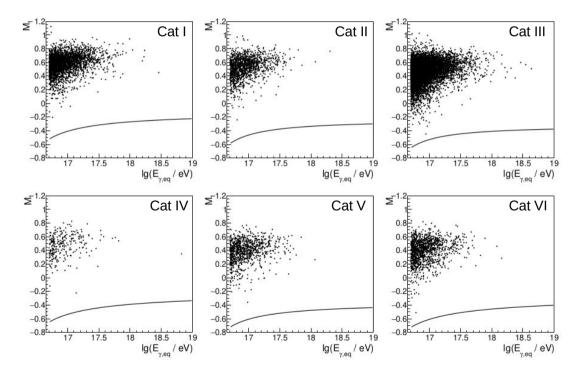
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- ▶ Since M₁ scales with the number of muon density measurements, any missing UMD station would bias M₁
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- ► M₁^{cut} is the "photon candidate cut"
- ► It is threshold value at which a 50% signal efficiency is reached
- ► M₁^{cut} scales with the event muon footprint, i.e., with the primary energy and the number of UMD stations

Full data-set unblinding: no photon candidates

► Events for each of the six categories. Lines correspond to the parametrized M₁^{cut} (the "photon candidate cuts")



- ► No events below the photon candidate cuts
- \rightarrow upper limits to the photon flux, Φ_{y} , above an energy $E_{y}^{\,th}$

$$\Phi_{\gamma} < \frac{3.095}{(1 - f_{\text{burnt}}) \times f_{\text{cut},\gamma} \times \epsilon_{\gamma}}$$

- ► f_{burnt} = 0.1 (fraction of burnt data)
- ► $f_{cut,y} = 0.504$ (signal efficiency)
- \triangleright ϵ_v , exposure to photons

E_{γ}^{th} (PeV)	$\epsilon_{\gamma} \; (\mathrm{km^2 sr yr})$	$\Phi_{\gamma} \; (\mathrm{km}^{-2} \mathrm{sr}^{-1} \mathrm{yr}^{-1})$	$\Phi_{\gamma} + \sigma_{\rm sys} \; ({\rm km}^{-2} {\rm sr}^{-1} {\rm yr}^{-1})$
50	(0.58 ± 0.02)	12.3	13.8
80	(0.61 ± 0.03)	11.7	13.5
120	(0.63 ± 0.03)	11.3	13.3
200	(0.63 ± 0.03)	11.3	13.6

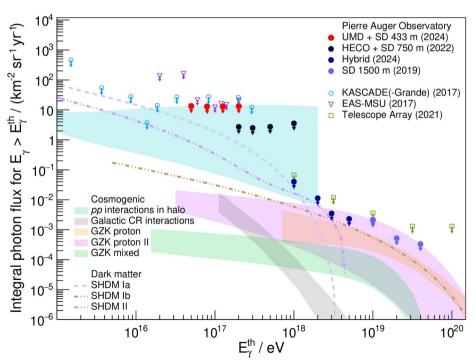
► Upper limits increased by:

- energy bias of the photon-equivalent scale (12-17)% - variations in the spectral index (0.8-2.2)%

▶ No impact from muon deficit or UMD simulation

Upper limits to the diffuse photon flux

- ► First search for a diffuse flux of primary photons above between 50 and 200 PeV from the Southern Hemisphere
- ► Opportunity to constrain the mass-lifetime phase-space for specific super-heavy dark matter models and to explore the photon flux from proton-proton interactions in the Galactic halo
- ► Groundwork for a nearly real-time search for primary photons in the tens of PeV domain



Search for a diffuse flux of photons with energies above tens of PeV at the Pierre Auger Observatory

The Pierre Auger Collaboration
E-mail: auger spokespersons@fnal.gov

Abstract. Diffuse photons of energy above 0.1 PeV, produced through the interactions between cosmic rays and either interstellar matter or background radiation fields, are powerful tracers of the distribution of cosmic rays in the Galaxy. Furthermore, the measurement of a diffuse photon flux would be an important probe to test models of super-heavy dark matter decaying into gamma-rays. In this work, we search for a diffuse photon flux in the energy range between 50 PeV and 200 PeV using data from the Pierre Auger Observatory. For the first time, we combine the air-shower measurements from a 2m² surface array consisting of 19 water-Cherenkov surface detectors, spaced at 433 m, with the muon measurements from a an array of buried scintillators placed in the same area. Using 15 months of data, collected while the array was still under construction, we derive upper limits to the integral photon flux ranging from 13.3 to 13.8 km⁻² sr⁻¹ yr⁻¹ above tens of PeV. We extend the Pierre Auger Observatory photon search program towards lower energies, covering more than three decades of cosmic-ray energy. This work lays the foundation for future diffuse photon searches: with the data from the next 10 years of operation of the Observatory, this limit is expected to improve by a factor of ~ 20.

- ► Currently, the Auger photon search programme covers three orders of magnitude in energy
- ► Searches at higher energies exploit the air-shower longitudinal development and integrated signals and timing observables from the larger SD arrays (see backup)

JCAP 05 (2025) 061

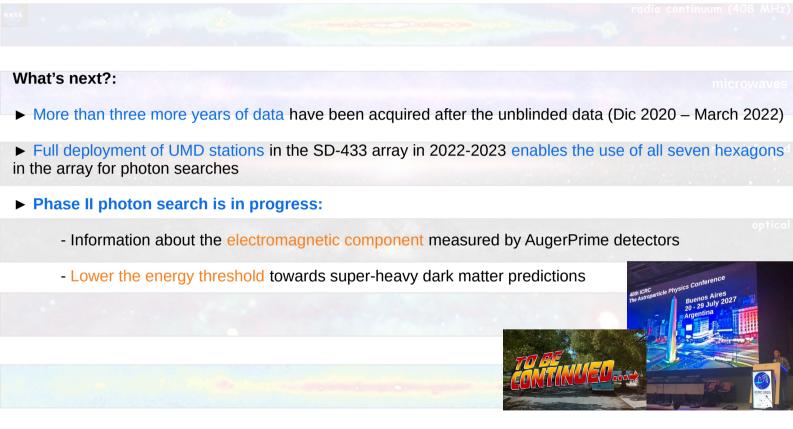
A key ingredient: the DDAp programme





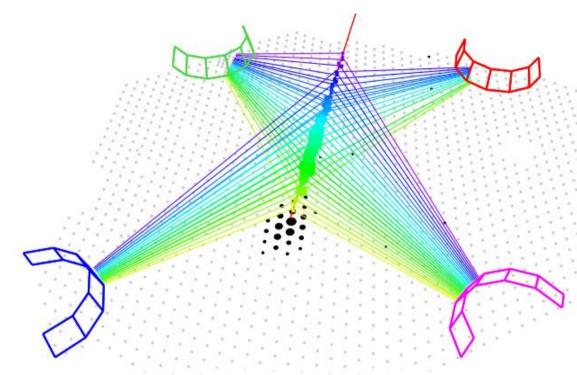
- ► Very important regarding the academic and professional development
- ► Unique opportunity to explore a variety of cultures and get to know many wonderful persons!

Unveiling a new face of the Milky Way

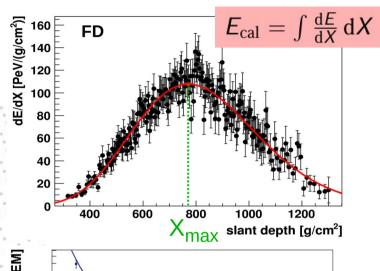


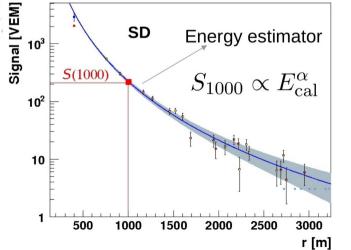
Backup

Hybrid detection of air showers



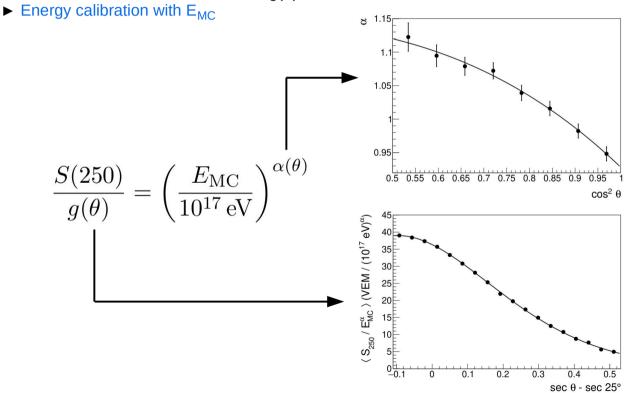
- ▶ Primary energy from the FD measurements
- \blacktriangleright SD energy estimator calibrated against FD \rightarrow SD-only reconstruction
- ▶ Maximum of shower development, X_{max}, is sensitive to the mass composition

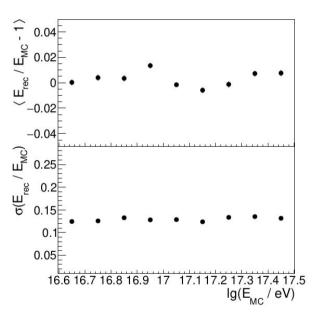




Photon energy reconstruction

- \blacktriangleright NKG with one slope parameter (β) used to reconstruct the LDF
- ightharpoonup Dedicated parametrization of β for photon primaries
- ▶ Gaisser-Hillas attenuation function $g(\theta)$





► <2% bias, 12% resolution

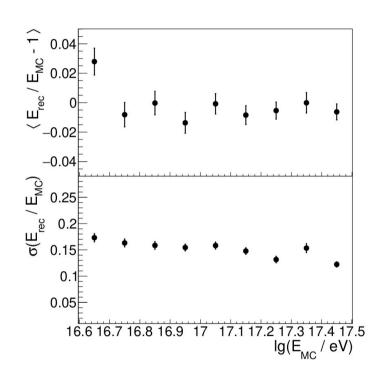
Proton energy reconstruction

- ▶ LDF model developed with Phase I data
- Attenuation function g'(θ) resembling CIC
- ▶ Direct energy calibration: unbinned likelihood to estimate the five free parameters

$$\frac{S(250)}{g'(\theta)} = \left(\frac{E_{\rm MC}}{10^{17} \,\text{eV}}\right)^{\underline{\alpha'}}$$

$$g'(\theta) = \underline{a'} \times \left(1 + \underline{b'} \times x + \underline{c'} \times x^2 + \underline{d'} \times x^3\right)$$

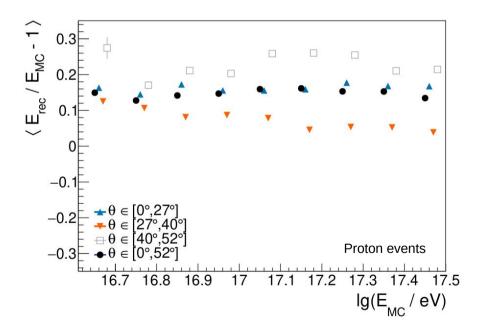
$$x = \cos^2 \theta - \cos^2 30^{\circ}$$



► <2% bias, 18% to 13% resolution

Photon-equivalent energy scale

► A unified energy scale is essential for accurately comparing events initiated by different primary species



- ► If reconstructed with the photon scale, proton events have a 15% bias (due to muonic component in hadronic showers)
- ► Noticeable angular dependence → photon and proton showers attenuate differently in the atmosphere
- → idea: let's combine the photon and proton energy calibrations!

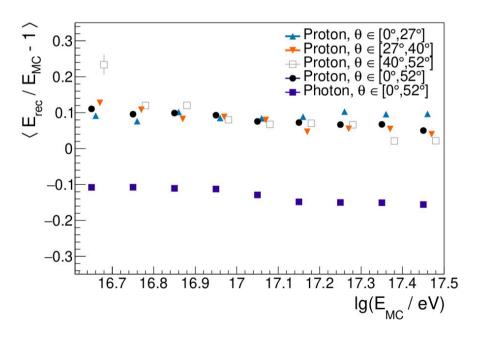
$$\frac{S(250)}{g'(\theta)} = \left(\frac{E_{\rm MC}}{10^{17} \, \text{eV}}\right)^{\alpha(\theta)}$$



Proton

Photon-equivalent energy scale

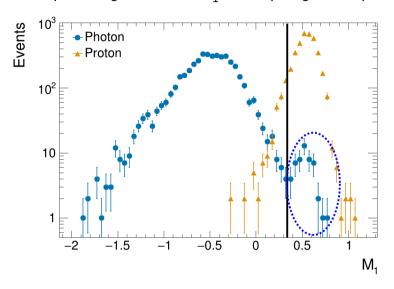
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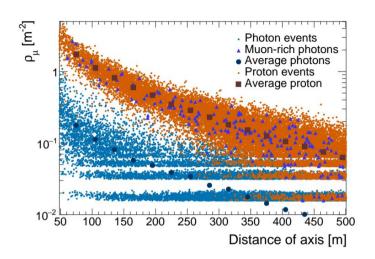
- ► Proton events reconstructed with the photon-equivalent scale are assigned energies between 5% and 10% higher than E_{MC}
- ► Photon events reconstructed with the photon-equivalent scale are assigned energies between 10% and 15% lower than E_{MC}
- ightharpoonup In a given $E_{\gamma,eq}$ bin, protons with lower true energy coexist with photons with higher true energy
 - → a conservative scenario for discrimination

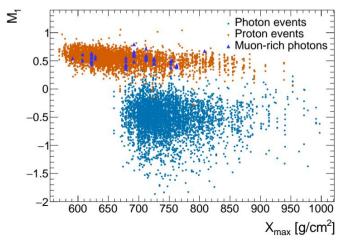
Muon-rich photon showers between 10^{16.7} eV and 10^{16.8} eV

► Muon-rich photon showers are defined as those producing events with $M_1 > 0.35$ (~ bkg median)

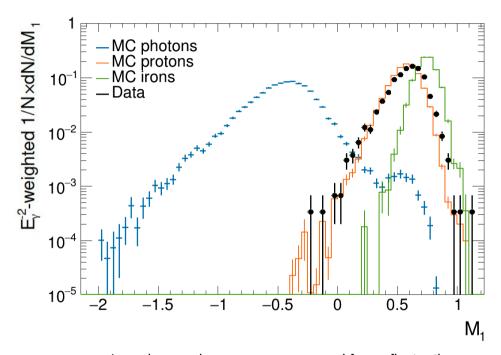


- ► Muon-rich photon showers have a muon content and a longitudinal development compatible with an average proton shower
- \blacktriangleright These special showers have a leading π^{\pm} at the first stages of the development

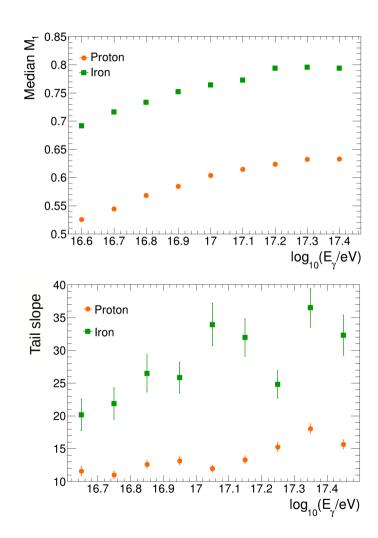




M₁ iron distributions



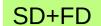
- $\,\blacktriangleright\,$ Iron showers have more muons and fewer fluctuations
- → larger M₁ and steeper tail slopes
- ▶ Both features lead to smaller bkg contamination
- ▶ Note that iron MC set size is 25% of the protons one

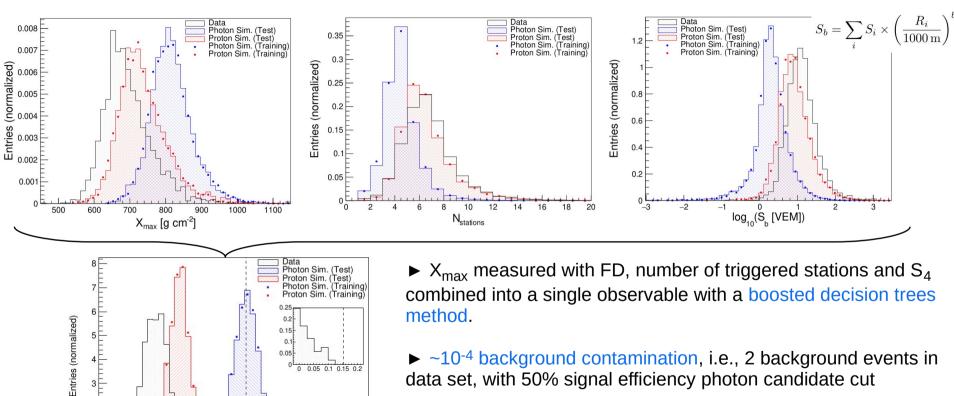


Photon search at $2x10^{17}$ eV $- 10^{18}$ eV

-0.2

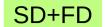
Boosted decision trees observable

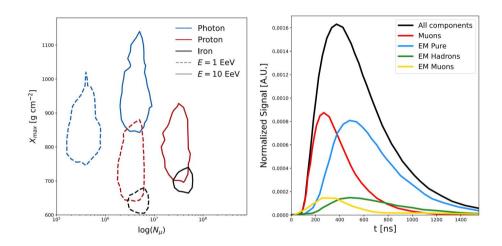




- ▶ ~10⁻⁴ background contamination, i.e., 2 background events in data set, with 50% signal efficiency photon candidate cut
- Data set spanning 5.5 years → ~2.5 km² sr yr
- ▶ 0 photon candidate events

Photon search at 10^{18} eV – 10^{19} eV

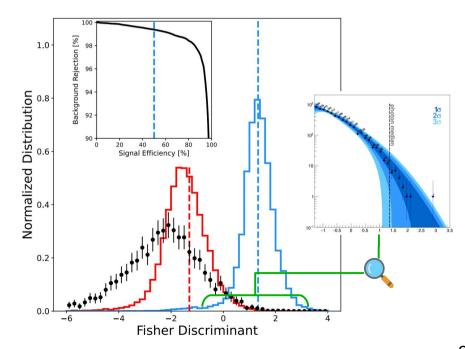




- $ightharpoonup X_{max}$ measured with the FD. F_{μ} is a proxy of the muon content extracted from the SD signals
- ightharpoonup Estimated by matching the measured signal in SD stations to S_{pred} which is a decomposition in EM and muonic components ($S_{i,comp}$)

$$S_{\text{pred}} = \sum_{i=1}^{4} S_i = \sum_{i=1}^{4} f_i(F_{\mu}) S_{i,\text{comp}}$$

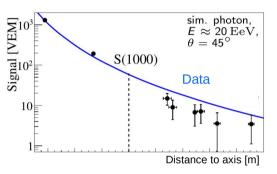
- Data set spans 12 years → ~1,000 km² sr yr
- ▶ 22 photon candidate events, but compatible with background of 30 ± 15 events estimated with data

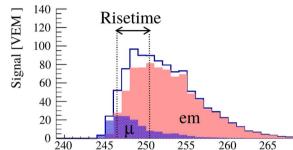


Photon search at $> 10^{19} \text{ eV}$









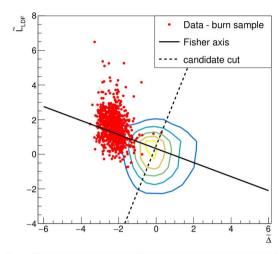
Longer signal rise-time

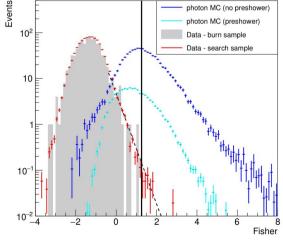
$$L_{\text{LDF}} = \log_{10} \left(\frac{1}{N} \sum_{i=1}^{N} \frac{S_i}{S_{\text{LDF}}(r_i)} \right)$$

$$\Delta = \frac{1}{N} \sum_{i=1}^{N} \frac{t_{1/2}^{i} - t_{1/2}^{\text{bench}}}{\sigma_{t_{1/2}}^{i}}$$

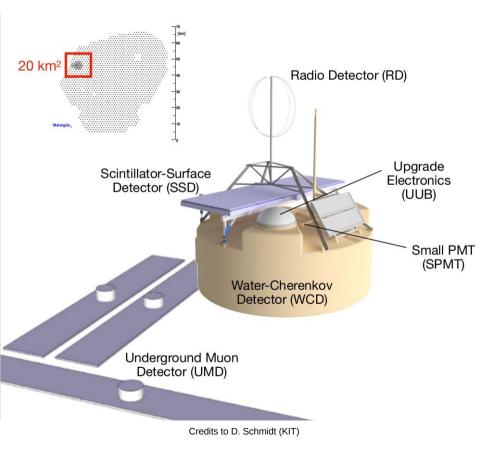
Time [25 ns]

- ► Benchmarks t_{1/2}^{bench} and S_{LDF} obtained with data
- → method free of cosmic-ray composition assumptions
- ► Fisher analysis trained with photon simulations and a subset of data
- ► Data set spans 16 years → 17,000 km² sr yr
- ► 16 photon candidate events, but compatible with the background contamination





The AugerPrime Upgrade



- ▶ Data can be exploited to (re-)define discriminating observables:
- e / μ decomposition using UMD/SSD (θ < 60°)
- shower development through radio measurements ($\theta > 60^{\circ}$)
- measured signals closer to the core with the SPMT
- ► Observables can be combined using cutting-edge algorithms (neural networks, AI, etc)
- ► Extending mass sensitivity to the lower energies (< 10¹⁷ eV) and improving photon/hadron discrimination

